Geometry and gradient flow structure of the arctangential heat flow

Yann Brenier, CNRS, DMA/ENS, 45 rue d'Ulm, FR-75005 Paris, in association with the CNRS-INRIA team "MOKAPLAN".

Gradient flows: challenges and new directions, ICMS, 10 to 14 September 2018, Edinburgh.

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Outline

We introduce the "arctangential" heat flow

 $\partial_t D = \Delta(\arctan D)$



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We introduce the "arctangential" heat flow

 $\partial_t D = \Delta(\arctan D)$

and show that this degenerate scalar parabolic equation has a hidden gradient flow structure, in parallel with the mean curvature flow for graphs

$$\partial_t \phi = \sqrt{1 + |\nabla \phi|^2} \ \nabla \cdot \left(\frac{\nabla \phi}{\sqrt{1 + |\nabla \phi|^2}} \right)$$

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The arctangential heat flow

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also reads, in non conservative form for $D = tan(\pi \psi)$,

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In (saturated) discrete form and 2 space dimensions:

$$\frac{\psi_{i,j}^{n+1} - \psi_{i,j}^{n}}{\cos(\pi\psi_{i,j}^{n})^{2}} = \frac{\psi_{i+1,j}^{n} + \psi_{i-1,j}^{n} + \psi_{i,j+1}^{n} + \psi_{i,j-1}^{n}}{4} - \psi_{i,j}^{n}.$$

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Few numerical results.

The 1d arctangential heat flow: brownian initial condition, $x \in \mathbb{R}/\mathbb{Z}$.

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Numerical solution for 256 grid points, 4096 time steps.

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The 2d arctangential heat flow: level sets of the data, $x \in \mathbb{R}^2/\mathbb{Z}^2$.

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The 2d arctangential heat flow: noisy initial condition, $x \in \mathbb{R}^2/\mathbb{Z}^2$.

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The 2d arctangential heat flow: recovery of the level sets $x \in \mathbb{R}/\mathbb{Z}, t \ge 0.$

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The 2d arctangential heat flow: other choice of data (with binary values), $x \in \mathbb{R}^2/\mathbb{Z}^2$.

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The 2d arctangential heat flow: noisy initial condition, $x \in \mathbb{R}^2/\mathbb{Z}^2$.

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The 2d arctangential heat flow: recovery of the level sets $x \in \mathbb{R}^2/\mathbb{Z}^2$, $t \ge 0$.

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twitter.com <u>Gabriel Peyré on Twitter</u> <u>"Yann Brenier just introduced and studied the arctangential heat equation, related to Born-Infeld</u> theory of Electromagnetism. This equation quantizes the values of the functions. https://t.co /NDdtl1QC9C https://t.co/kQ1W0uHfwn"

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The main result (Y.B. 2018, HAL preprint: hal-01740320): The equation of extremal graphs in Minkowski's space

$$\partial_t (\frac{\partial_t \phi}{R}) = \nabla \cdot (\frac{\nabla \phi}{R}), \quad R = \sqrt{1 - \partial_t \phi^2 + |\nabla \phi|^2}$$

(which is a nonlinear wave equation) generates two twin gradient flows.

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(which is a nonlinear wave equation) generates two twin gradient flows. The first one is the arctangential heat flow $\partial_t D = \Delta(\arctan D)$, while the second one is just the well known mean curvature flow for graphs

$$\partial_t \phi = \sqrt{1 + |\nabla \phi|^2} \ \nabla \cdot \left(\frac{\nabla \phi}{\sqrt{1 + |\nabla \phi|^2}} \right)$$

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The arctangential flow $\partial_t D = \lambda \Delta(\arctan(D\lambda^{-1}))$ can be easily written in optimal transport style (à la Otto):

 $\partial_t D = \nabla \cdot (D \ \nabla(\mathcal{F}'(D))),$

The arctangential flow $\partial_t D = \lambda \Delta(\arctan(D\lambda^{-1}) \operatorname{can} \operatorname{be} \operatorname{easily} \operatorname{written} \operatorname{in} \operatorname{optimal} \operatorname{transport} \operatorname{style}$ (à la Otto):

 $\partial_t D = \nabla \cdot (D \ \nabla(\mathcal{F}'(D))), \text{ where}$ $\mathcal{F}(D) = D \log \left(\frac{D}{\sqrt{1+D^2\lambda^{-2}}}\right) - \lambda \arctan(D\lambda^{-1})$

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 $\mathcal{F}(D) = D \log \left(\frac{D}{\sqrt{1+D^2\lambda^{-2}}}\right) - \lambda \arctan(D\lambda^{-1}) \text{ is the}$ Legendre transform of $u \to \lambda \arcsin(\lambda^{-1}e^u)$ (extended by $+\infty$ for $u > \log \lambda$), which can be seen as a "catastrophic" version of the usual exponential.

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The "catastrophic" exponential for different values of parameter λ . $u \rightarrow \lambda \arcsin(\lambda^{-1} \exp(u))$ (extended by $+\infty$ for $u > \log \lambda$).

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The inverse of the "catastrophic" exponential (after symmetrization and periodization), for different values of λ $\nu \rightarrow \frac{1}{2} \log(\lambda^2 \sin^2(\nu \lambda^{-1}))$

is also used in "optimal unbalanced transport theory"!

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Gradient flow structure of the arctangential flow

Alternately (and more geometrically), the arctangential flow is the gradient flow of functional $D \rightarrow \int_{\mathbb{T}^d} \sqrt{1 + D^2}$ with metric $||v||_D^2 = \int_{\mathbb{T}^d} (1 + D^2)^{-1/2} |vD|^2$, tangent vectors at point *D* being written as $\dot{D} = -\nabla \cdot (vD)$.

Gradient flow structure of the arctangential flow

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$$\partial_t (\frac{\partial_t \phi}{R}) = \nabla \cdot (\frac{\nabla \phi}{R}), \quad R = \sqrt{1 - \partial_t \phi^2 + |\nabla \phi|^2}.$$

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Gradient flow structure of the MC flow for graphs

In a parallel way, the mean curvature flow for graphs can be interpreted as the gradient flow of functional $B \rightarrow \int_{\mathbb{T}^d} \sqrt{1 + |B|^2}$ with metric

 $||v||_B^2 = \int_{\mathbb{T}^d} (1 + |B|^2)^{-1/2} (B \cdot v)^2$ where tangent vectors at point $B = \nabla \phi$ are written as $\dot{B} = -\nabla (B \cdot v)$.

Gradient flow structure of the MC flow for graphs

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 $||v||_B^2 = \int_{\mathbb{T}^d} (1 + |B|^2)^{-1/2} (B \cdot v)^2$ where tangent vectors at point $B = \nabla \phi$ are written as $\dot{B} = -\nabla (B \cdot v)$. Again, this structure comes from the nonlinear wave equation

$$\partial_t (\frac{\partial_t \phi}{R}) = \nabla \cdot (\frac{\nabla \phi}{R}), \quad R = \sqrt{1 - \partial_t \phi^2 + |\nabla \phi|^2}.$$

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Final part: proof of our main result.

We want to derive from the nonlinear wave equation

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Final part: proof of our main result.

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$$\partial_t (\frac{\partial_t \phi}{R}) = \nabla \cdot (\frac{\nabla \phi}{R}), \quad R = \sqrt{1 - \partial_t \phi^2 + |\nabla \phi|^2},$$

at once, both the arctangential heat flow

 $\partial_t D = \Delta(\arctan D)$

and the mean curvature flow for graphs

$$\partial_t \phi = \sqrt{1 + |\nabla \phi|^2} \ \nabla \cdot \left(\frac{\nabla \phi}{\sqrt{1 + |\nabla \phi|^2}} \right)$$

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Proof/First step. As $\phi(t, x)$ solves the equation of extremal surfaces in Minkowski's space, then

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$$(D, B, P) = \frac{1}{\sqrt{1 - \partial_t \phi^2 + |\nabla \phi|^2}} (\partial_t \phi, \nabla \phi, -\partial_t \phi \nabla \phi)$$

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$$(D, B, P) = \frac{1}{\sqrt{1 - \partial_t \phi^2 + |\nabla \phi|^2}} (\partial_t \phi, \nabla \phi, -\partial_t \phi \nabla \phi)$$

solves the "entropic" system of conservation laws:

$$\partial_t B + \nabla \left(\frac{P \cdot B - D}{h}\right) = 0, \quad \partial_t D + \nabla \cdot \left(\frac{PD - B}{h}\right) = 0,$$
$$\partial_t P + \nabla \cdot \left(\frac{P \otimes P + B \otimes B}{h}\right) = \nabla \left(\frac{1 + B^2}{h}\right),$$
with $h = \sqrt{1 + D^2 + B^2 + P^2}$ as convex entropy.

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Proof/Second step. To get a gradient flow, we follow Y.B. and Xianglong Duan (ARMA 2018) by performing the simple singular change of time $t \rightarrow \theta = t^2/2$:

$$\mathcal{B}(heta, x) = \mathcal{B}(\sqrt{2 heta}, x),$$

 $\mathcal{D}(heta, x) = rac{\mathcal{D}(\sqrt{2 heta}, x)}{\sqrt{2 heta}}, \quad \mathcal{P}(heta, x) = rac{\mathcal{P}(\sqrt{2 heta}, x)}{\sqrt{2 heta}},$

requiring initial condition D = P = 0 at t = 0(which corresponds to $\partial_t \phi(0, x) = 0$ in terms of ϕ).

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In a somewhat dual way, a second natural change is

$$\mathcal{D}(\theta, x) = \mathcal{D}(\sqrt{2\theta}, x),$$
$$\mathcal{B}(\theta, x) = \frac{\mathcal{B}(\sqrt{2\theta}, x)}{\sqrt{2\theta}}, \quad \mathcal{P}(\theta, x) = \frac{\mathcal{P}(\sqrt{2\theta}, x)}{\sqrt{2\theta}},$$
requiring initial condition $\mathcal{B} = \mathcal{P} = 0$ at $t = 0$ (which corresponds to $\nabla \phi = 0$ at $t = 0$ for ϕ).

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After performing the change of time $t \rightarrow \theta = t^2/2$, we get, in the 1st case, the non automous system:

$$egin{aligned} \partial_{ heta}\mathcal{B} &=
abla \left(\mathcal{D} - rac{\mathcal{P}\cdot\mathcal{B}}{\mathcal{H}}
ight), \ \mathcal{H} &= \sqrt{\mathbf{1} + \mathcal{B}^2 + 2 heta(\mathcal{D}^2 + \mathcal{P}^2)}, \ \mathcal{D} -
abla \cdot \left(rac{\mathcal{B}}{\mathcal{H}}
ight) &= -2 heta \left(\partial_{ heta}\mathcal{D} +
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abla \left(rac{1 + \mathcal{B}^2}{\mathcal{H}}
ight),$

$$\partial_{\theta}\mathcal{B} = \nabla\left(\mathcal{D} - \frac{\mathcal{P}\cdot\mathcal{B}}{\mathcal{H}}
ight), \ \mathcal{H} = \sqrt{1+\mathcal{B}^2}$$

$$\mathcal{D} = \nabla \cdot \left(rac{\mathcal{B}}{\mathcal{H}}
ight), \ \ \mathcal{P} + \nabla \cdot \left(rac{\mathcal{B} \otimes \mathcal{B}}{\mathcal{H}}
ight) = \nabla \left(rac{1 + \mathcal{B}^2}{\mathcal{H}}
ight),$$

Symmetrically, the second rescaling leads to the arctangential heat equation

$$\partial_{\theta} \mathcal{B} = \nabla \left(\mathcal{D} - \frac{\mathcal{P} \cdot \mathcal{B}}{\mathcal{H}} \right), \ \mathcal{H} = \sqrt{1 + \mathcal{B}^2}$$

$$\mathcal{D} =
abla \cdot \left(rac{\mathcal{B}}{\mathcal{H}}
ight), \ \ \mathcal{P} +
abla \cdot \left(rac{\mathcal{B} \otimes \mathcal{B}}{\mathcal{H}}
ight) =
abla \left(rac{1 + \mathcal{B}^2}{\mathcal{H}}
ight),$$

Symmetrically, the second rescaling leads to the arctangential heat equation and, then, the twin gradient flow structures easily follow.

$$\partial_{\theta}\mathcal{B} = \nabla\left(\mathcal{D} - \frac{\mathcal{P}\cdot\mathcal{B}}{\mathcal{H}}\right), \ \mathcal{H} = \sqrt{1+\mathcal{B}^2}$$

$$\mathcal{D} =
abla \cdot \left(rac{\mathcal{B}}{\mathcal{H}}
ight), \ \ \mathcal{P} +
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ight),$$

Symmetrically, the second rescaling leads to the arctangential heat equation and, then, the twin gradient flow structures easily follow. End of proof.

$$\partial_{\theta}\mathcal{B} = \nabla\left(\mathcal{D} - \frac{\mathcal{P}\cdot\mathcal{B}}{\mathcal{H}}
ight), \ \mathcal{H} = \sqrt{1+\mathcal{B}^2}$$

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abla \cdot \left(rac{\mathcal{B}}{\mathcal{H}}
ight), \ \ \mathcal{P} +
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Symmetrically, the second rescaling leads to the arctangential heat equation and, then, the twin gradient flow structures easily follow. End of proof. **THANKS !**

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