Riemannian Geometry for Computational Anatomy

Introducing the LDDMM framework.

Jean Feydy October 10, 2017

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Jean Feydy (2016-2019) :

- PhD student under the supervision of Alain Trouvé.
- Caïman at the ENS.

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Notes for this talk are available online (in French) :

www.math.ens.fr/~feydy/Teaching/

- Culture Mathématique, chap. 4-6.
- Introduction à la Géométrie Riemannienne par l'Étude des Espaces de Formes.

Introduction



Research in Image Processing :

• Signal analysis.

Figure 1: Image denoising, from [2]. .



Figure 1: Brain segmentation, from [7]. .

Research in Image Processing :

- Signal analysis.
- Segmentation.



Figure 1: Brain database, from [4].

Research in Image Processing :

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Figure 1: Brain database, from [4].

Research in Image Processing :

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We need appropriate representations.









JPEG2000, JPEG : Wavelets, Blockwise (high + low) frequencies



(a) Original image. (b) JPEG2000, 20 : 1. (c) JPEG, 20 : 1.

Figure 3: Taken from www.photozone.de.



Figure 4: CNN visualization, from vision03.csail.mit.edu/cnn_art/.



Figure 5: Reference image.



Figure 5: With a transferred texture component. [6]



Figure 5: With a transferred texture component. [6]



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Convolutional Neural Networks : Texture-invariant Classification



Figure 6: CNNs allow tech companies to group together photos and sketches of beavers.

How do we handle intra-class variability?



Figure 7: Silhouettes segmented from a fishing net. [3]

Rigid Body Analysis

From images to labeled point clouds



Figure 8: Anatomical landmarks on a tuna fish. [1]

Let $X, Y \in \mathbb{R}^{M \times D}$ be two labeled point clouds. Let $S_{\tau, \upsilon}$ denote the **rigid**-body transformation of parameters τ (translation) and υ (rotation + scaling). Then, try to find

$$\begin{aligned} \pi_0, \upsilon_0 &= \arg\min_{\tau, \upsilon} \qquad \|S_{\tau, \upsilon}(\mathbf{X}) - \mathbf{Y}\|_2^2 \tag{1} \\ &= \arg\min_{\tau, \upsilon} \sum_{m=1}^M |\upsilon \cdot \mathbf{x}^m + \tau - \mathbf{y}^m|^2. \end{aligned}$$

Position, Scale and Orientation



Figure 9: Matching the blue wing on the red one. (Wikipedia)

Pros and cons of Rigid body analysis

Pros:

- Simple and robust
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This model is a standard pre-processing tool. However, it is too limited to allow in-detail analysis. **Optimal Transport**

Let: (x^1, \ldots, x^M) , (y^1, \ldots, y^M) be two point clouds in \mathbb{R}^D .

Find a collection of paths $\gamma^m : t \in [0, 1] \mapsto \gamma^m_t$, a permutation $\sigma : [1, M] \to [1, M]$ such that

 $\forall m, \quad \gamma_0^m = x^m \text{ and } \gamma_1^m = y^{\sigma(m)}, \tag{3}$

minimizing

$$\ell^{2}(\gamma) = \sum_{m=1}^{M} \int_{t=0}^{1} \|\dot{\gamma}_{t}^{m}\|^{2} \,\mathrm{d}t.$$
(4)

$$X \xrightarrow{\gamma} Y.$$
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$$X \xrightarrow{\gamma} Y.$$
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If we relabel the unit masses (x^1, \ldots, x^M) and (y^1, \ldots, y^M) , find a **permutation** $\sigma : \llbracket 1, M \rrbracket \to \llbracket 1, M \rrbracket$ minimizing

$$C^{\mathbf{X},\mathbf{Y}}(\sigma) = \sum_{m=1}^{M} \left\| \mathbf{x}^m - \mathbf{y}^{\sigma(m)} \right\|^2.$$
 (6)

 σ is an optimal labeling.

Image matching as a mass-carrying problem



Figure 10: Optimal transport between two curves seen as mass distributions : from a déblai to a remblai.
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Pros and cons of Optimal Transport

Pros:

- Well-posed, convex problem
- Global and precise matchings
- Light-speed numerical solvers at hand (Cuturi, 2013)

Cons:

• Discards topology : tears shapes apart

This model is mathematically and numerically appealing. However, it does not provide any <mark>smoothness</mark> guarantee.

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This model is mathematically and numerically appealing. However, it does not provide any smoothness guarantee. Can we build a rich and practical model for smooth deformations ?

The LDDMM framework









The LDDMM framework

Regularized transport : a Riemannian problem

A naive way to regularize transport : Find $\sigma : \llbracket 1, M \rrbracket \to \llbracket 1, M \rrbracket$ minimizing $C_{k}^{\mathbf{X}, \mathbf{Y}}(\sigma) = \underbrace{\sum_{m} \left\| \mathbf{x}^{m} - \mathbf{y}^{\sigma(m)} \right\|^{2}}_{\text{Displacement cost}} + \underbrace{\sum_{m,m'} k(\mathbf{x}^{m}, \mathbf{x}^{m'}) \cdot \left\| \mathbf{y}^{\sigma(m)} - \mathbf{y}^{\sigma(m')} \right\|^{2}}_{\text{Regularization cost}},$ (7)

with k(x, y) a kernel neighborhood function.

Find a permutation $\sigma: \llbracket 1, M \rrbracket \to \llbracket 1, M \rrbracket$ minimizing

$$C_{k,\text{sym}}^{X,Y}(\sigma) = \underbrace{\sum_{m} \left\| x^m - y^{\sigma(m)} \right\|^2}_{\text{Displacement cost}} + \frac{1}{2} \underbrace{\sum_{m,m'} k(x^m, x^{m'}) \cdot \left\| y^{\sigma(m)} - y^{\sigma(m')} \right\|^2}_{X \to Y \text{ regularization cost}} + \frac{1}{2} \underbrace{\sum_{m,m'} k(y^m, y^{m'}) \cdot \left\| x^{\sigma^{-1}(m)} - x^{\sigma^{-1}(m')} \right\|^2}_{Y \to X \text{ regularization cost}}$$

This cost is symmetric, but does not handle properly the shapes between X and Y.

Find a collection of paths γ^m from **X** to **Y** minimizing

$$C_{k}(\gamma) = \int_{0}^{1} \left[\underbrace{\sum_{m} \|\dot{\gamma}_{t}^{m}\|^{2}}_{\text{Displacement cost}} + \underbrace{\sum_{m,m'} k(\gamma_{t}^{m}, \gamma_{t}^{m'}) \cdot \left\|\dot{\gamma}_{t}^{m} - \dot{\gamma}_{t}^{m'}\right\|^{2}}_{\text{Regularization cost}} \right] \text{d}t.$$

Particles will move optimally if they are :

- lazy
- gregarious wrt. their k-neighbors

With $\gamma_t = (\gamma_t^1, \dots, \gamma_t^M) \in \mathbb{R}^{M \times D}$, we can write

$$C_k(\gamma) = \int_0^1 \dot{\gamma}_t^{\mathsf{T}} g_{\gamma_t} \dot{\gamma}_t \mathrm{d}t.$$
 (8)

Optimal deformations are geodesics on the space of landmarks $\mathbb{R}^{M \times D}$ endowed with a Riemannian metric g_q :

$$\frac{\left(\mathrm{d}_{g}(q \to q + v \cdot \mathrm{d}t)\right)^{2}}{\mathrm{d}t} = \sum_{m} \|v^{m}\|^{2} + \sum_{m,m'} k(q^{m}, q^{m'}) \cdot \left\|v^{m} - v^{m'}\right\|^{2}$$
$$= v^{\mathrm{T}}g_{q}v = \|v\|_{g_{q}}^{2}$$
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The LDDMM framework

Geodesic shooting on a Riemannian manifold

Riemann : conveniently working with arbitrary geometries



(a) As a deformed square.

(b) Embedded in \mathbb{R}^3 .

Figure 12: The donut-shaped torus.

Sometimes, we can compute geodesics explicitly...



Figure 13: Explicit geodesics on homogeneous manifolds.
(b) is adapted from www.pitt.edu/~jdnorton/.

But this is not the case in general



Figure 14: Geodesics on the Duhem's bull, embedded in \mathbb{R}^3 . Taken from www.chaos-math.org. Geodesic \Longrightarrow locally "straight" \Longrightarrow second order ODE, the **geodesic equation** satisfied by $\gamma_t = (\gamma_t^1, \dots, \gamma_t^D)$:

$$\forall d \in \llbracket 1, D \rrbracket, \quad \dot{\gamma}_t^d = -\sum_{1 \leqslant i, j \leqslant D} \Gamma_{jj}^d(\gamma_t) \cdot \dot{\gamma}_t^i \dot{\gamma}_t^j, \tag{10}$$

where the Christoffel symbols $\Gamma^d_{ij}(\gamma_t)$ are given by :

$$\Gamma_{ij}^{d}(\gamma_{t}) = \frac{1}{2} \sum_{l=1}^{D} g^{dl}(q) \cdot \left(\partial_{i}g_{jl}(q) + \partial_{j}g_{il}(q) - \partial_{l}g_{ij}(q)\right), \quad (11)$$

with g_{ij} the metric tensor and g^{dl} its inverse, the cometric.

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The "Christoffel" equation is an ODE on the tangent bundle :

$$(q_t, v_t) = (\gamma_t, \dot{\gamma}_t). \tag{12}$$

Hamilton : one should work on the cotangent bundle :

$$(q_t, p_t) = (q_t, g_{q_t} v_t). \tag{13}$$

We denote $\mathsf{K}_q=g_q^{-1}$ and $\mathsf{H}(q,p)=rac{1}{2}p^\mathsf{T}\mathsf{K}_q p$, so that

$$\frac{1}{2} v_t^{\mathsf{T}} g_{q_t} v_t = \underbrace{\frac{1}{2} \|\dot{\gamma}_t\|_{\gamma_t}^2}_{\text{Kinetic energy}} = \frac{1}{2} p_t^{\mathsf{T}} \mathcal{K}_{q_t} p_t = \mathcal{H}(q_t, p_t).$$
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Hamilton, 1833

 γ_t is a geodesic if and only if the lifted cotangent trajectory (q_t, p_t) follows the Hamiltonian equation :

$$\begin{cases} \dot{q}_t = +\frac{\partial H}{\partial p}(q_t, p_t) = +K_{q_t}p_t \\ \dot{p}_t = -\frac{\partial H}{\partial q}(q_t, p_t) = -\frac{1}{2}\partial_q(p_t, K_q p_t)(q_t) \end{cases}$$
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In the cotangent phase space, we flow along the symplectic gradient :

$$X(q,p) = \begin{pmatrix} +\frac{\partial H}{\partial p}(q,p) \\ -\frac{\partial H}{\partial q}(q,p) \end{pmatrix} = "R_{-90}\circ"(\nabla H(q,p)).$$
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Consider a free-falling particle of mass m:

$$q = z,$$
 $v = \dot{z},$ (17)
 $\dot{q} = v,$ $\dot{v} = -g.$ (18)

Now, we can write $\mathbf{p} = \mathbf{m}\mathbf{v}$ so that

$$H(q,p) = "E_{cin}"(q,p) + "E_{pp}"(q,p) = \frac{1}{2}\frac{p^2}{m} + mgq.$$
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$$\begin{cases} \dot{q} = +\frac{\partial H}{\partial p} = +p/m \\ \dot{p} = -\frac{\partial H}{\partial q} = -mg \end{cases}.$$
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A geodesic path γ_t is characterized by (q_0, p_0) .

To compute any geodesic starting from a source q_0 , we simply need a shooting momentum p_0 and a simplistic Euler scheme :

$$\begin{cases} q_{t+0.1} = q_t + 0.1 \cdot K_{q_t} p_t \\ p_{t+0.1} = p_t - 0.1 \cdot \partial_q (p_t, K_q p_t) (q_t) \end{cases}$$
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Exponential map :

$$\operatorname{Exp}_{q_0}: p_0 \in T^*_{q_0} \mathcal{M} \mapsto q_1 \in \mathcal{M}$$
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We are looking for :

- Tearing-adverse metrics on the space of landmarks
- Efficient ways to compute geodesics (deformations)

Hamilton has taught us that :

- Geodesics are "simple" iff the cometric $K_q = g_q^{-1}$ is simple
- The Exponential map can be computed efficiently

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Kernel cometrics and Diffeomorphic trajectories

Parallelism is the way forward



Figure 15: Highly-parallel MoKaMachine (Mokaplan Inria team).

GPUs in action



Figure 16: *Mythbusters Demo GPU versus CPU*, from the Nvidia YouTube channel.

Use a reduced kernel matrix

$$k_{q} = \begin{pmatrix} k(q^{1}, q^{1}) & k(q^{1}, q^{2}) & \cdots & k(q^{1}, q^{M}) \\ k(q^{2}, q^{1}) & k(q^{2}, q^{2}) & \cdots & k(q^{2}, q^{M}) \\ \vdots & \vdots & \ddots & \vdots \\ k(q^{M}, q^{1}) & k(q^{M}, q^{2}) & \cdots & k(q^{M}, q^{M}) \end{pmatrix}$$
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so that

$$H(q,p) = \frac{1}{2}p^{\mathsf{T}} K_{q} p = \frac{1}{2} \sum_{i,j=1}^{M} k(q^{i},q^{j}) \cdot p^{i\mathsf{T}} p^{j}.$$
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In a computational sense, this is the simplest family of cometrics on the space of points clouds.

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Let k be a smooth enough kernel function, with $\widehat{k}(\omega) \in \mathbb{R}_+^*$. If $v : \mathbb{R}^D \to \mathbb{R}^D$ is a vector field on the ambient space, define

$$\|\boldsymbol{v}\|_{k}^{2} = \int_{\boldsymbol{\omega}\in\mathbb{R}^{D}} \frac{1}{\widehat{k}(\boldsymbol{\omega})} |\widehat{\boldsymbol{v}}(\boldsymbol{\omega})|^{2} d\boldsymbol{\omega}.$$
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- + $\mathcal{V}_k = \{ v \mid \|v\|_k < \infty \}$ is a Hilbert space of k-smooth vector fields
- We assume k is smooth enough, so that $\delta_x : v \mapsto v(x)$ belongs to the dual space $(U_k)^*$: we link with the theory of **Reproducing** Kernel Hilbert Spaces.

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Assume that (v_t) is a time-varying vector field such that

$$\ell_k(v)^2 = \int_0^1 \|v_t\|_k^2 \, \mathrm{d}t < \infty.$$
 (26)

According to Picard-Lindelöf theorem, we can integrate the flow, find a unique trajectory φ_t of diffeomorphisms such that for every point $x \in \mathbb{R}^D$ and time $t \in [0, 1]$:

$$\begin{split} \varphi_0(\mathbf{x}) &= \mathbf{x} \qquad \text{and} \quad \frac{\mathrm{d}}{\mathrm{d}t} \left[\varphi_t(\mathbf{x}) \right] = \mathbf{v}_t \circ \varphi_t(\mathbf{x}), \\ \text{i.e.} \quad \varphi_0 &= \operatorname{Id}_{\mathbb{R}^D} \quad \text{and} \qquad \varphi_t &= \operatorname{Id}_{\mathbb{R}^D} + \int_{s=0}^t \mathbf{v}_s \circ \varphi_s \, \mathrm{d}s. \end{split}$$

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According to Picard-Lindelöf theorem, we can integrate the flow, find a unique trajectory φ_t of diffeomorphisms such that for every point $x \in \mathbb{R}^D$ and time $t \in [0, 1]$:

$$\begin{split} \varphi_0(\mathbf{x}) &= \mathbf{x} \qquad \text{and} \quad \frac{\mathrm{d}}{\mathrm{d}t} \left[\varphi_t(\mathbf{x}) \right] = \mathbf{v}_t \circ \varphi_t(\mathbf{x}), \\ \text{i.e.} \quad \varphi_0 &= \mathrm{Id}_{\mathbb{R}^D} \quad \text{and} \qquad \varphi_t = \mathrm{Id}_{\mathbb{R}^D} + \int_{s=0}^t \mathbf{v}_s \circ \varphi_s \, \mathrm{d}s. \end{split}$$

We define $G_k = \{\varphi_1 \mid \cdots\}$ the set of diffeomorphisms obtained by integrating *finite-cost* vector flows $(v_t) \in L^2(U_k)$.

 G_k is an infinite-dimensional Riemannian manifold modeled on U_k . As diffeomorphisms carry around images and measures, we try to minimize

$$C^{2}(\varphi_{1}) = \ell_{k}(v)^{2} = \int_{0}^{1} ||v_{t}||_{k}^{2} dt < \infty$$
 (27)

under the constraint that

$$X \xrightarrow{\varphi_1} Y.$$
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$$C^{2}(\varphi_{1}) = \ell_{k}(\nu)^{2} = \int_{0}^{1} \|\nu_{t}\|_{k}^{2} dt < \infty$$
 (27)

under the constraint that

$$\mathbf{X} \xrightarrow{\varphi_1} \mathbf{Y}.$$
 (28)

Reduction Principle

Let q_t be a time-dependent point cloud, k a kernel function. Then, the two propositions below are equivalent :

i) q_t is a geodesic for the kernel cometric K_q , with momentum p_t associated to the Hamiltonian

$$H(q,p) = \frac{1}{2}p^{\mathsf{T}} \mathcal{K}_q p. \tag{29}$$

ii) q_t is carried around by a locally optimal diffeomorphic trajectory $\varphi_t = \text{Flow}(v_t)$, and we have

$$v_t = k \star p_t$$
 i.e. $v_t(x) = \sum_{m=1}^{M} k(q_t^m, x) p_t^m$. (30)

At any time t,

$$v_t = \arg\min\{\|v\|_k \mid \forall \, m, v(q_t^m) = v_t(q_t^m)\}.$$
(31)

Hence, as v_t does not have any superfluous component,

$$v_t \in \{ v \mid \forall m, v(q_t^m) = 0 \}^{\perp_k}$$
(32)

i.e.
$$v_t \in \left(\bigcap_{m=1}^{M} \left\{ v \mid \left\langle \delta_{\mathbf{q}_t^m}, v \right\rangle = 0 \right\} \right)^{\perp_k}$$
. (33)

But we also know that :

At any time t,

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Hence, as v_t does not have any superfluous component,

$$\nu_t \in \{ \nu \mid \forall m, \nu(q_t^m) = 0 \}^{\perp_k}$$
(32)

i.e.
$$v_t \in \left(\bigcap_{m=1}^{M} \left\{ v \mid \left\langle \delta_{\mathbf{q}_t^m}, v \right\rangle = 0 \right\} \right)^{\perp_k}$$
. (33)

But we also know that :

$$\left\langle \mathbf{k} \star \delta_{q_{t}^{m}}, \boldsymbol{\nu} \right\rangle_{\mathbf{k}} = \int_{\boldsymbol{\omega} \in \mathbb{R}^{D}} \frac{1}{\widehat{k}(\boldsymbol{\omega})} \overline{\widehat{k} \star \delta_{q_{t}^{m}}(\boldsymbol{\omega})} \cdot \widehat{\boldsymbol{\nu}}(\boldsymbol{\omega}) \, \mathrm{d}\boldsymbol{\omega}$$
(34)

$$= \int_{\omega \in \mathbb{R}^{D}} \overline{\widehat{\delta_{q_{t}^{m}}(\omega)}} \cdot \widehat{\nu}(\omega) \, \mathrm{d}\omega$$
(35)

$$= \left\langle \delta_{q_t^m}, \nu \right\rangle = \nu(q_t^m). \tag{36}$$

Hence why, at any time t,

$$\begin{aligned}
\nu_{t} \in \left(\bigcap_{m=1}^{M} \left\{ \left. \nu \right| \left\langle k \star \delta_{q_{t}^{m}}, \nu \right\rangle_{k} = 0 \right\} \right)^{\perp_{k}} \\
= \bigcup_{m=1}^{M} \left(k \star \delta_{q_{t}^{m}} \right)^{\perp_{k} \perp_{k}}
\end{aligned} \tag{37}$$

$$= \operatorname{Vect}\left(k \star \delta_{q_t^m}, \ m \in \llbracket 1, M \rrbracket\right).$$
(39)

So, one can write

$$\nu_t = k \star \left(\sum_{m=1}^M p_t^m \, \delta_{q_t^m} \right) = k \star p_t, \tag{40}$$

and

$$\|\boldsymbol{v}_t\|_k^2 = \left\langle k \star p_t, k^{(-1)} \star k \star p_t \right\rangle = \left\langle k \star p_t, p_t \right\rangle = \boldsymbol{p}_t^\mathsf{T} \boldsymbol{K}_{\boldsymbol{q}_t} \boldsymbol{p}_t.$$
(41)

Hence why, at any time t,

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(41)
























































































We have now presented the Large Deformation Diffeomorphic Metric Mapping, or LDDMM setting :

- OT $(\sigma = 0) \xrightarrow{\sigma + +} G_k \xrightarrow{\sigma + +} (\sigma = +\infty)$ Translations
- Deformations computed through geodesic shooting

- Hamilton's theorem $(q_a \longrightarrow K_a)$
- The current availability of GPUs
- The Reduction Principle

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- OT $(\sigma = 0) \xrightarrow{\sigma + +} G_k \xrightarrow{\sigma + +} (\sigma = +\infty)$ Translations
- Deformations computed through geodesic shooting

The (basic) framework relies on three pillars :

- Hamilton's theorem $(g_q \longrightarrow K_q)$
- The current availability of GPUs (parallelism)
- The Reduction Principle

 $((q_t, p_t) \longleftrightarrow \varphi_t)$
Conclusion

We can now emulate D'Arcy Thompson's work

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which fossils are subject (as we have seen on p. 811) as the result of shearing-stresses in the solid rock.

Fig. 519 is an outline diagram of a typical Scaroid fish. Let us deform its rectilinear coordinates into a system of (approximately) coaxial circles, as in Fig. 520, and then filling into the new system,





Fig. 517. Argyropelecus Olferei.

Fig. 518. Sternoptyz diaphana.

space by space and point by point, our former diagram of Souray, we obtain a very good outline of an allied fish, belonging to a neighbouring family, of the genus *Domaonthus*. This case is all the more interesting, because upon the body of our *Pomaonthus* there are striking colour bands, which correspond in direction very closely



to the lines of our new curved ordinates. In like manner, the still more bicarce outlines of other finishes of the same family of Chaetodonts will be found to correspond to very slight modifications of similar coordinates; in other words, to small variations in the values of the constants of the coaxia curves.

In Figs. 521-524 I have represented another series of Acanthopterygian fishes, not very distantly related to the foregoing. If we

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start this series with the figure of *Polyprion*, in Fig. 521, we see that the outlines of *Pseudopriacanthus* (Fig. 522) and of *Sebastes* or *Scorpaena* (Fig. 523) are easily derived by substituting a system





Fig. 521. Polyprion

Fig. 522. Pseudopriacanthus altus.

of triangular, or radial, coordinates for the rectangular ones in which we had inscribed *Polyprion*. The very curious fish *Antigonia capros*, an oceanic relative of our own boar-fish, conforms, closely to the peculiar deformation represented in Fig. 524.



Fig. 525 is a common, typical *Diodos* or porcupine-fish, and in Fig. 526 I have deformed its vertical coordinates into a system of concentric circles, and its horizontal coordinates into a system of curves which, approximately and provisionally, are made to resemble

Figure 25: Excerpt from the seminal book of D'Arcy Wentworth Thompson (1860-1948), *On Growth and Forms*.

Biologists, Neurologists and Physicians would like to conduct statistical surveys such as :

- Linear regression
- Mean computation + Principal Component Analysis
- Transport of tangential information

Problem : no meaningful algebraic structure $(+, \times)$ on shapes.

Given a mere Riemannian distance, we provide :

- Geodesic regression
- Fréchet Mean + PCA on shooting momentums
- Parallel transport

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Transfer of anatomical data: animated silhouettes



Figure 26: Video presentation of the (non-LDDMM) paper *Anatomy Transfer Fast Forward*, Siggraph Asia 2013 by Ali-Hamadi, Liu, Gilles et al.

Transfer of anatomical data: medical applications



Figure 27: Video presentation of the (non-LDDMM) paper *Anatomy Transfer Fast Forward*, Siggraph Asia 2013 by Ali-Hamadi, Liu, Gilles et al.

Construction of anatomical atlases



Figure 28: Building an atlas from the retina dataset [5].

A continuum of professions



Figure 29: A (very) schematic view of the fields related to Computational Anatomy.

A continuum of professions



Figure 29: The people behind the labels.















Thank you for your attention.

The pytorch library: symbolic maths on the GPU

```
31
     def Hqp(q, p, sigma) :
32
          "The hamiltonian, or kinetic energy of the shape q with momentum p."
          pKqp = _k(q, q, sigma) * (p a p.t()) # pKqp_{i,i} = k(q_i, q_i) \langle p_i, p_i \rangle_2
33
34
          return .5 * pKap.sum()
                                                   # H(q, p) = \frac{1}{2} \sum_{i \in i} k(q_i, q_j) p_i p_i
35
36
     # The partial derivatives of the Hamiltonian are automatically computed !
     def dq Hqp(q.p.sigma) :
37
38
          return torch.autograd.grad( Hgp(g.p.sigma), g. create graph=True)[0]
     def dp Hqp(q,p,sigma) :
39
40
          return torch.autograd.grad( Hqp(q,p,sigma), p, create graph=True)[0]
41
42
     def HamiltonianShooting(q, p, sigma) :
43
          "Shoots to time 1 a k-geodesic starting (at time 0) from g with momentum p."
          for t in range(10) :
                                                      # Let's hardcode the "dt = .1".
44
              q,p = [q + .1 * _dp_Hqp(q,p,sigma) , # Euler steps for the Hamiltonian flow
45
                      p - .1 * dq Hqp(q,p,sigma) ] # in the cotangent bundle.
46
          return [q.p]
                                                      # Return the final state + momentum.
47
```

Shooting routines emulate the Exponential map



A deformation $\varphi_t(X)$ is encoded as a shooting momentum $p_0 \in T^*_X \mathcal{M}$.

Find the momentum $X \xrightarrow{\varphi \simeq p_0} Y$ through gradient descent.

Figure: Matching a curve to another.

Shooting routines emulate the Exponential map



A deformation $\varphi_t(X)$ is encoded as a shooting momentum $p_0 \in T^*_X \mathcal{M}$.

Find the momentum $X \xrightarrow{\varphi \simeq p_0} Y$ through gradient descent.

Figure: Matching a curve to another.



Figure 31: Iteration 0.



Figure 31: Iteration 3.



Figure 31: Iteration 4.



Figure 31: Iteration 5.



Figure 31: Iteration 6.



Figure 31: Iteration 7.



Figure 31: Iteration 8.



Figure 31: Iteration 9.



Figure 31: Iteration 10.



Figure 31: Iteration 11.



Figure 31: Iteration 12.



Figure 31: Iteration 13.



Figure 31: Iteration 14.



Figure 31: Iteration 15.



Figure 31: Iteration 16.


Figure 31: Iteration 17.



Figure 31: Iteration 18.



Figure 31: Iteration 19.



Figure 31: Iteration 20.



Figure 31: Iteration 21.



Figure 31: Iteration 22.



Figure 31: Iteration 23.



Figure 31: Iteration 24.



Figure 31: Iteration 25.



Figure 31: Iteration 26.



Figure 31: Iteration 27.



Figure 31: Iteration 28.



Figure 31: Iteration 29.



Figure 31: Iteration 30.



Figure 31: Iteration 31.



Figure 31: Iteration 32.



Figure 31: Iteration 33.



Figure 31: Iteration 34.



Figure 31: Iteration 35.



Figure 31: Iteration 36.



Figure 31: Iteration 37.



Figure 31: Iteration 38.



Figure 31: Iteration 39.



Figure 31: Iteration 41.



Figure 31: Iteration 42.



Figure 31: Iteration 43.



Figure 31: Iteration 44.



Figure 31: Iteration 46.



Figure 31: Iteration 47.



Figure 31: Iteration 48.



Figure 31: Iteration 49.



Figure 31: Iteration 50.



Figure 31: Iteration 52.



Figure 31: Iteration 53.



Figure 31: Iteration 54.



Figure 31: Iteration 55.


Figure 31: Iteration 56.



Figure 31: Iteration 57.



Figure 31: Iteration 58.



Figure 31: Iteration 59.



Figure 31: Iteration 60.



Figure 31: Iteration 61.



Figure 31: Iteration 62.



Figure 31: Iteration 64.



Figure 31: Iteration 65.



Figure 31: Iteration 66.



Figure 31: Iteration 67.



Figure 31: Iteration 68.



Figure 31: Iteration 69.



Figure 31: Iteration 70.



Figure 31: Iteration 71.



Figure 31: Iteration 72.



Figure 31: Iteration 73.



Figure 31: Iteration 74.



Figure 31: Iteration 75.



Figure 31: Iteration 77.



Figure 31: Iteration 78.



Figure 31: Iteration 79.



Figure 31: Iteration 80.



Figure 31: Iteration 81.



Figure 31: Iteration 82.



Figure 31: Iteration 83.



Figure 31: Iteration 85.



Figure 31: Iteration 86.



Figure 31: Iteration 87.



Figure 31: Iteration 88.



Figure 31: Iteration 89.



Figure 31: Iteration 90.

Matchings of partially observed shapes



(a) X and Y.



(b) Target Y, view 1.



(c) Target Y, view 2.







(d) Source X. (e) *f*(*X*), view 1. (f) *f*(*X*), view 2.

Figure 32: Matching artifacts for the retina dataset.

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